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# A DESORBED GAS MOLECULAR IONIZATION MECHANISM FOR ARCING ONSET IN SOLAR ARRAYS IMMERSSED IN A LOW-DENSITY PLASMA

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## Abstract

Previous experimental studies have hypothesized that the onset of Solar Array Arc (SAA) initiation in low-density space plasmas is caused by a desorbed gas molecular ionization mechanism. Indeed past investigations performed at the NASA Glenn Plasma Interaction Facility tend to not only support the desorbed gas molecular ionization mechanism, but have gone as far as identifying the crucial molecular species that must be present for molecular ion dominated process to occur. When electrical breakdown occurs at a triple junction site on a solar array panel, a quasi-neutral plasma cloud is ejected. Assuming the main component of the expelled plasma cloud by weight is due to water vapor, the fastest process available is due to

H<sub>2</sub>O molecules and OH<sup>+</sup> ions, or more succinctly, dissociative molecular-ion dominated recombination processes: H<sub>2</sub>O<sup>+</sup> + e<sup>-</sup> → H\* + OH\*. Recently published spectroscopic observations of solar array arc spectra in ground tests have revealed the well-known molecular OH band (302nm-309nm), as well as the molecular SiH band (387nm peak) and, the molecular CH band (432nm peak). Note that the OH band is observed in emission arcs where water vapor is present. Strong atomic lines were also observed for H<sub>β</sub> at 486nm and H<sub>α</sub> at 656.3nm in prior ground testing. Independent supporting evidence of desorbed gas molecular ionization mechanisms also come from measurements of arc current pulse widths at different capacitances. We will revisit an earlier first order approximation demonstrating the dependence of arc current pulse widths on the square root of the capacitance. The simple arc current pulse width model will be

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then be used to estimate the temperature of the arc plasma (currently believed to be somewhere in the range of 3 to 5 eV). The current paper then seeks to extend the outlined work by including numerous vacuum chamber measurements obtained with a quadrupole mass spectrometer. A small solar array was mounted inside the vacuum chamber. A plasma source, also mounted inside the vacuum chamber, is used to simulate a low-density plasma environment. The solar array is then biased to a high negative potential and allowed to arc while a mass spectrometer is used to record the partial pressure of H<sub>2</sub>O and to track other significant changes in mass (1-150) AMU.

### Nomenclature

$C$	= capacitance, $\mu F$
$V_s$	= spacecraft potential, V
$V_p$	= plasma potential, V
$V_o$	= array output voltage, V
$U$	= bias potential, V
$N_e$	= electron number density, $cm^{-3}$
$T_e$	= electron temperature, eV
$A$	= atomic molecular mass, AMU
$m_e$	= electron rest mass, kg
$k$	= Boltzmann constant, $J K^{-1}$
$e$	= electron charge, q
$\gamma$	= recombination rate, $cm^3 s^{-1}$
$y, x$	= curve fitted variables
$a, b$	= intercept, slope constants
$\tau$	= pulse width, s

### Introduction

The trend for modern day satellite designers is to provide robust high voltage power generation equipment that is both lightweight and arc resistant in the harsh plasma environment of space. Solar arrays are often employed to generate the required power because they fit most of the design criteria set forth above. While older payloads typically operated at less than 80 V, newer payloads are designed to operate at 300V or more. The high

voltage requirements have placed a heavier demand on modern spacecraft power systems.

Traditionally spacecraft designers have incorporated a single point ground to connect the negative end of the solar array directly to the structure while the entire spacecraft is allowed to electrically float in the space plasma environment. The negative grounding scheme has led to potentially hazardous problems involving arcing [1-5].

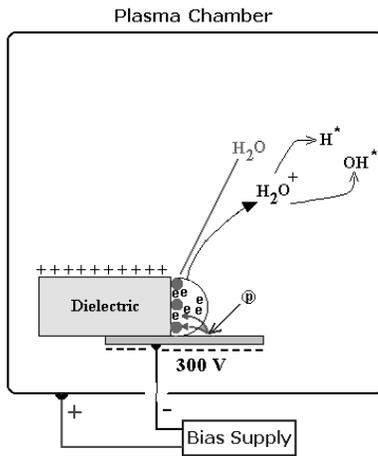
If the un-powered spacecraft is allowed to float in the ambient plasma it will float at a potential  $V_s$  close to the plasma potential  $V_p$ . Assume for the moment that the spacecraft structure floats at an approximate potential  $V_p = -2eV$ . Now add a negatively grounded array to the spacecraft structure and assume the array develops an output voltage  $V_o = -300V$ . The spacecraft potential is effectively pulled to a potential  $V_s \approx -70V$  [22].

Because the negative end of the solar array is grounded to the spacecraft structure, the metallic contacts of the array will acquire the same spacecraft potential  $V_s$  as the structure. The high negative potential  $V_s$  of the metallic contacts gives rise to a proportionate ion flux to the surface of the solar arrays, causing the individual solar cell cover glass slides to acquire a positive surface charge. Positively charged atoms accelerated towards the metallic contacts, will hit the contact, and release secondary electrons.

Previous experimental investigations have shown that arcs tend occur at the triple junction sites of a solar array. A triple junction site on a solar array is defined as any node where adjacent conductor- dielectric and plasma junctions intersect. The ongoing surface chemistry, due to the adsorption of water molecules by the metallic contacts coupled with secondary electron emission to the side surface of the dielectric, leads to electron impact desorption.

Electron Impact Desorption (EID) is believed to be responsible for ionization of

H<sub>2</sub>O molecules and the formation of H\* and OH\* [6-12]. Secondary electrons impacting a water vapor molecule causes the molecule to lose an electron and to acquire an overall positive charge: (H<sub>2</sub>O - e<sup>-</sup> → H<sub>2</sub>O<sup>+</sup>) Further bombardment of H<sub>2</sub>O<sup>+</sup> molecules by other secondary electron impacts leads to EID and causes dissociative recombination processes to emerge: (H<sub>2</sub>O<sup>+</sup> + e<sup>-</sup> → H\* + OH\*). See Figure 1 for details. Note the time scale for



**Figure (1): Electron impact ionization is believed to be the leading cause for emergence of dissociative recombination processes.**

dissociative recombination processes to occur is of the order of (8-10) $\mu$ s and the distances involved typically range from (1-2) cm. in radius from the triple junction site. Secondary electron emissions also cause the side surface of the cover glass to charge positive (see Figure1). The plasma sheath and mirror side charge are believed to be responsible for the formation of a strong E-field at the triple junction site. Increasing numbers of negative charges continue to build up on the side surface of the cover glass until an ensuing electron avalanche causes the arc to strike.

The scientific community has actively studied the problem of arc initiation and numerous models have been developed to

describe the underlying cause of arcing onset in space based solar arrays [13-21]. The so-called triple junction model has previously been employed (with limited success) to explain how the formation of a strong electric field at the metal, dielectric and plasma junctions of a solar array leads to arcing [13-16]. Past work performed at the Glenn Research Center (GRC) lends weight to our hypothesis that condensed water at the triple junction site is responsible for the desorbed gas ionization mechanism and the resulting dissociative molecular-ion recombination processes. The current work seeks to confirm the desorbed molecular gas hypothesis through an independent experimental verification method. The independent verification method used a quadruple mass spectrometer to track changes in the partial pressures of H<sub>2</sub>O, OH and H<sub>2</sub> as well as tracking changes to other mass species during the arcing process.

**Energy Estimates of SAA Plasmas:**

Direct evidence supporting a desorbed gas ionization mechanism has come from measurements of arc current pulse widths at different capacitances [16]. The dependence of arc current pulse widths on the square root of the capacitance can be written to a first order approximation as

$$(1) \quad t = \left( \frac{\gamma \cdot C \cdot U}{\pi \cdot e} \right)^{\frac{1}{2}} \cdot \left( \frac{A \cdot m_e}{2 \cdot k \cdot T_e} \right)^{\frac{3}{4}}$$

The dissociative recombination rate  $\gamma$  is defined as

$$(2) \quad \gamma = 10^{-7} \cdot \frac{1}{\sqrt{T_e}} \text{ cm}^3 \cdot \text{s}^{-1}$$

Where  $T_e$  is the electron temperature of the arc plasma in eV. Taking the atomic mass (A = 18) for the H<sub>2</sub>O molecule and a bias

potential  $U = 500V$  the current pulse width  $\tau$  in equation (1) can be rewritten as.

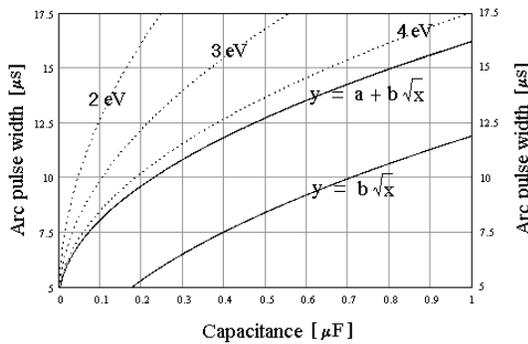
$$(3) \quad \tau = 53 \cdot C^{\frac{1}{2}} \cdot T_e^{-1} \quad \mu s$$

A simple square root approximation of capacitance to arc current pulses width data in equation (4) has yielded good results [16].

$$(4) \quad y = a + b \cdot \sqrt{x},$$

Arc Current pulse width vs. Capacitance Fitted to Square Root Approximation for Bias -500 V:

$$(y = a + b\sqrt{x}, a=4.31095, b=11.918383)$$



**Figure (2): Both solid lower curves were plotted with a value  $T_e = 4.45$  eV. All other dotted curves are plotted with the values of  $T_e$  indicated.**

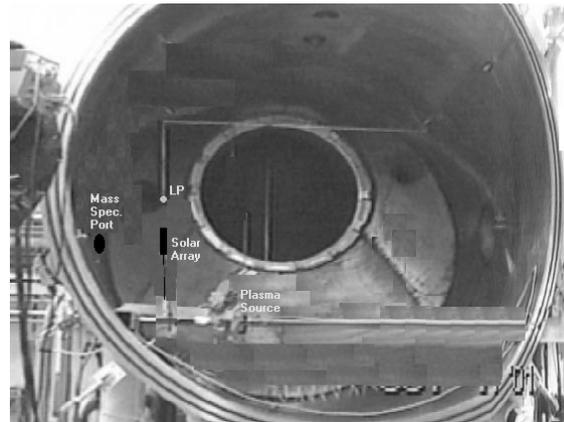
where:  $x$  represents the capacitance in  $\mu F$  and  $y$  the fitted arc pulse width in  $\mu s$ . The constants  $a$ ,  $b$  represent the respective y-intercept and slope (see Figure 2). Equations (3) & (4) can be used to estimate the energy  $T_e$  of the arc plasma by equating the following relations.

$$(5a) \quad 53 \cdot C^{\frac{1}{2}} \cdot T_e^{-1} = b \cdot \sqrt{x}$$

Assuming,  $x = C$  in equation (5a), we may solve directly for  $T_e$  to obtain

$$(5b) \quad T_e = \frac{53}{b} \cong 4.45 \quad eV$$

Figure 2 demonstrates the results obtained by plotting equation (4) at different electron temperatures (The capacitance,  $x$  is used as the running variable and the fitted arc pulse width  $y$  is the dependant variable). In Figure 2 the lower solid curve was plotted using equation (4) with the value  $a = 0$ ; Letting  $b = 53/T_e$  in equation (4) correctly results in equation (5a). The upper solid line plotted in Figure 2 was obtained using equation (4)

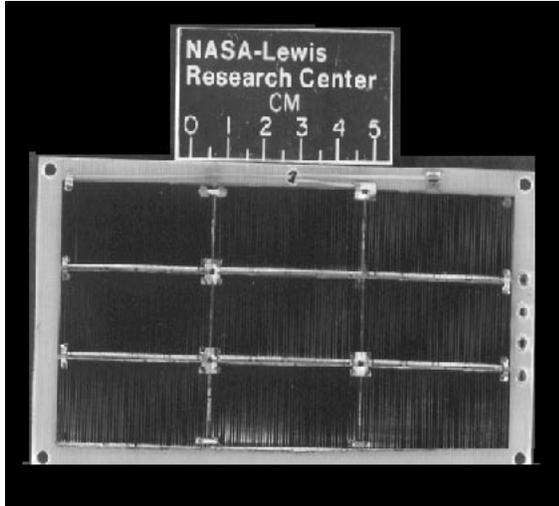


**Figure 3: Open view of Tenny vacuum chamber used in tests.**

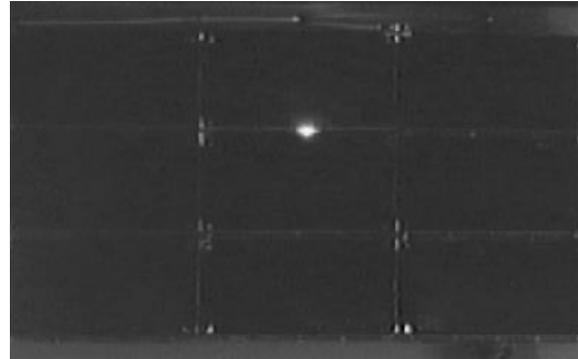
with the fitted values of both  $a$ , and  $b$  shown in Figure 2). It should also be noted that EID stimulated processes require far less energy (4.45 eV) to initiate than what is typically needed to ionize a hydrogen atom (13.6 eV).

### Experimental Apparatus

Ground tests for the mass species investigations were performed in the large (2.5m long by 1.8m diameter) Tenny vacuum chamber located at the NASA Glenn Plasma Interaction facility (See Figure 3). A commercially manufactured hollow cathode capable of producing  $T_e = 0.7$  eV electrons was used as a plasma source for the experiments. Xenon neutral gas for the experiments was carefully metered into the



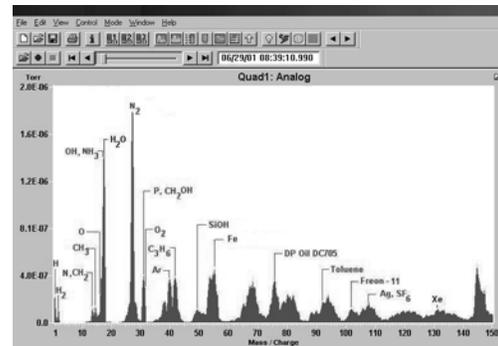
**Figure 4:** The pictured array is composed of 3 strings wired in parallel. Each string has 3 cells wired in series.



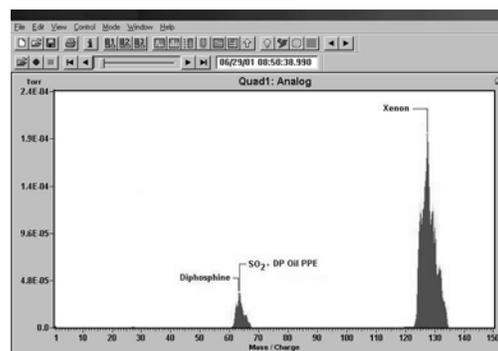
**Figure 5:** A typical arc on the solar array.  $C = 1 \mu\text{F}$ , bias voltage,  $U = 500 \text{ V}$ .

chamber using a precision vacuum valve with variable size orifice. An ionization gauge was then used to obtain the chamber pressure. The background pressure in the chamber was set at  $P_0 = 8 \times 10^{-5}$  Torr with the hollow cathode operating. The measured electron number densities ranged between  $N_e = (1 \times 10^4 - 8 \times 10^5)$  per  $\text{cm}^{-3}$  or about the same conditions used in the previous bell jar arc spectra experiments [17].

A small solar array test sample (Figure 4) was mounted inside the Teney vacuum chamber and allowed to float in the ambient plasma. The solar array contacts were shorted together and biased over a range of negative bias potentials,  $U = -(0.5 \text{ to } 1) \text{ kV}$  (See Figure 3). A range of negative potentials was used because the small exposed array contact area, coupled with sample conditioning effects, caused the arc rate to drop over time. The bias potential was continuously adjusted so that an arc rate of about 1 arc per second was achieved for the duration of the tests. Figure 5 features a typical arc on the biased solar array. A quadruple mass spectrometer was employed to monitor the change in partial



**Figure (6a)**



**Figure (6b)**

**Snapshots showing mass spectrum before Figure (6a) and after Figure (6b) xenon gas is introduced into the chamber.**

pressure of several mass species during the arcing process. The solar array sample was mounted a distance of 28cm from the feed-through port of the mass spectrometer (See Figure 3).

**Experimental Results**

Figure 6a shows a typical neutral gas background sweep for mass numbers (1-150)

AMU taken in the Tenny vacuum chamber just prior to adding xenon gas. Note the high concentrations of H<sub>2</sub>O and N<sub>2</sub> in the vacuum chamber. The concentrations of H<sub>2</sub>O and N<sub>2</sub> are approximately a factor of 5 larger in concentration than any other mass species recorded. Figure 6b was recorded just after establishing a constant

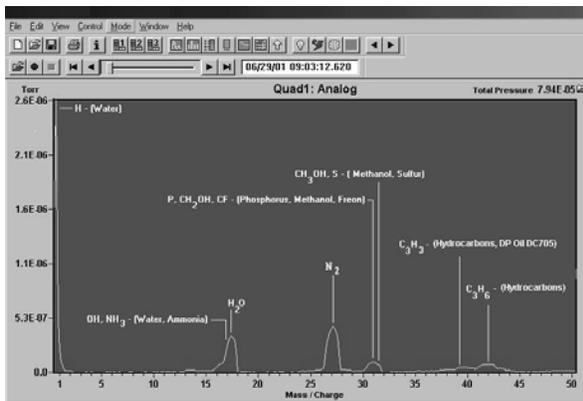


Figure (7a)

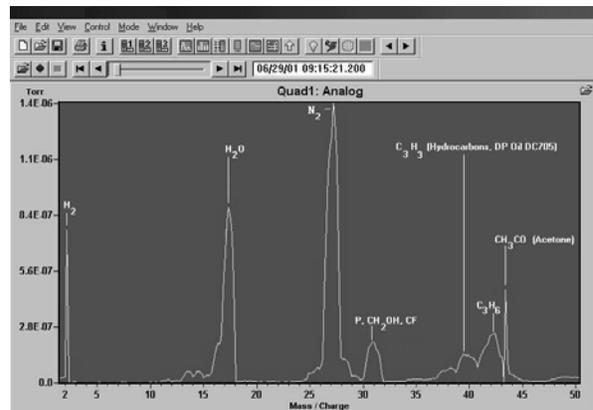


Figure (7c)

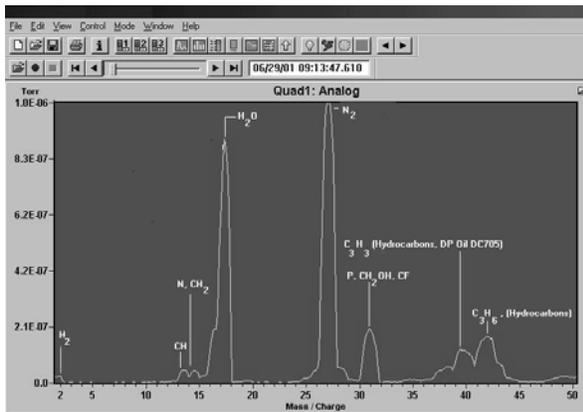


Figure (7b)

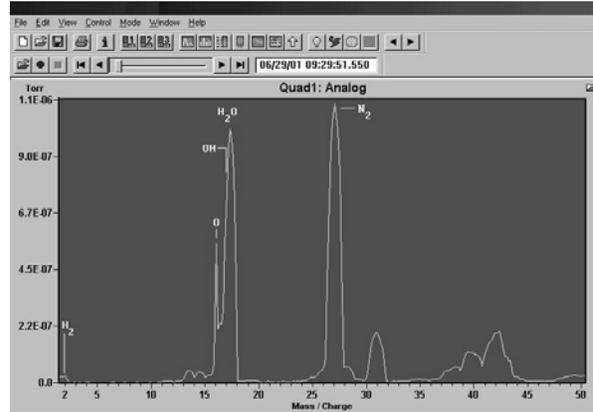


Figure (7d)

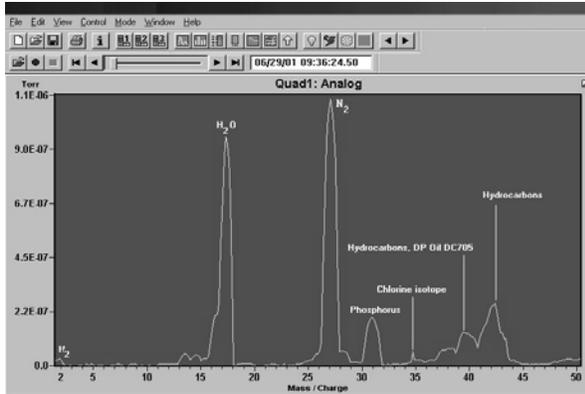


Figure (7e)

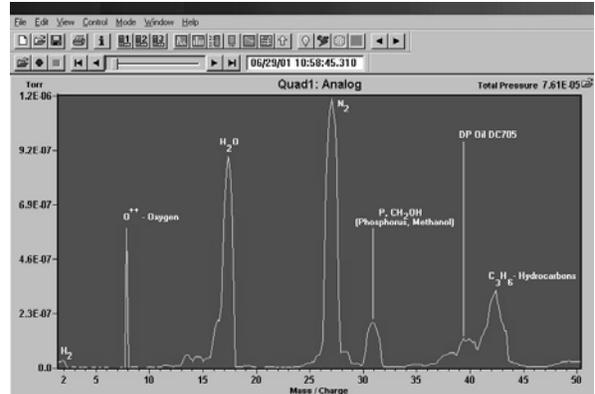


Figure (7g)

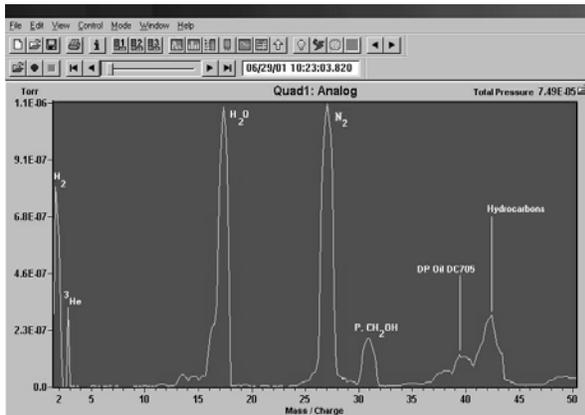


Figure (7f)

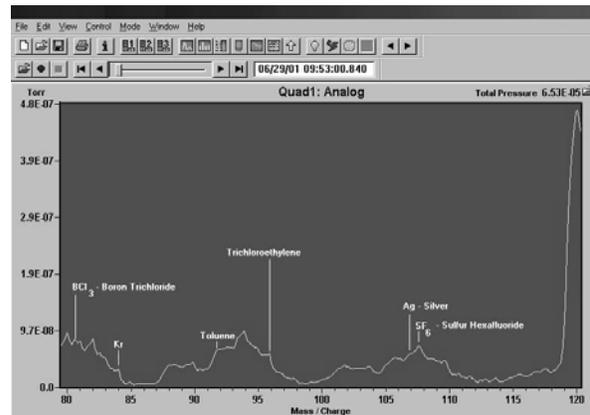


Figure (7h)

xenon flow rate for the hollow cathode. Note that the concentrations of the xenon gas peak and the  $\text{SO}_2$  peak are large and tend to drive the concentrations of all other mass species in the vacuum system to zero when plotted on this scale. In order to alleviate this problem it was decided to bracket the mass range for all successive sweeps (1-50 AMU and 80-120 AMU). Bracketing the mass range also seeks to cut the sweep interval time by less than one half (or just under 60 seconds for 1 complete sweep). Figures (7a-7g) are snapshots depicting the temporal evolution of the relative concentrations in partial pressure recorded for several mass spectra. The recorded data in Figures (7a-7h) were obtained with the plasma source operating and the solar array under a large negative potential bias.

Figure 7a was taken approximately 3.2 minutes after the array was biased. Note the large concentration of hydrogen and the smaller concentrations of  $\text{H}_2\text{O}$  and  $\text{N}_2$  in Figure 7a. The snapshot in Figure 7b was obtained approximately 10.6 minutes later. Note the increased levels of  $\text{H}_2\text{O}$  and  $\text{N}_2$  and the negligible  $\text{H}_2$  level in Figure 7b.

Figure 7c was obtained 1.6 minutes later still (or about 15.3 minutes into the experiment). Note the large spike corresponding to  $\text{H}_2$ . The  $\text{H}_2\text{O}$  concentration remains about the same level and there is a small increase in the partial pressure of  $\text{N}_2$  in Figure 7c. The  $\text{C}_3\text{H}_3$  and  $\text{C}_3\text{H}_6$  hydrocarbon concentrations continue to grow in figures 7a-7c but appears to level off after that.

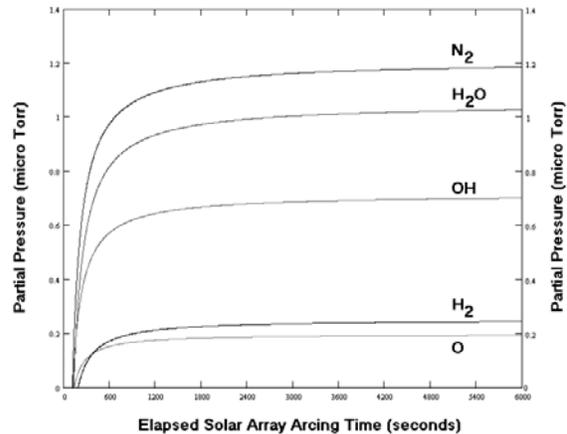
Figure 7d was taken 29.6 minutes into the experiment. Note the small concentration in the partial pressure of  $\text{H}_2$ . There appears to be

a large spike in the oxygen line (mass 16) in Figure 7d. Figure 7e was obtained 36.4 minutes into the experiment. The mass spectrum in Figure 7e appears to show the vacuum chamber has returned to its typical quiescent arcing state with no remarkable features present. Figure 7f was recorded 1.4 hours into the experiment. Note the large spikes observed for H<sub>2</sub> and the helium isotope, <sup>3</sup>He. Also note the similar levels in the concentrations of both H<sub>2</sub>O and N<sub>2</sub> lines recorded in figure 7f. The snapshot in Figure 7g demonstrates a doubly ionized oxygen atom, O III (O<sup>++</sup>, mass 8). Note the decreased level in the partial pressure of H<sub>2</sub>O and the large steady concentration of N<sub>2</sub> observed 1 hour and 58.75 minutes into the experiment. Finally the snapshot in Figure 7h depicts the typical mass spectrum observed for mass ranges (80-120) AMU. The mass species of the large peak recorded at 119 AMU is unknown.

Figure 8 illustrates the dependence of partial pressure on solar array arcing time for several mass species (N<sub>2</sub>, H<sub>2</sub>O, OH, H<sub>2</sub> and O) in the vacuum system. The plots were obtained by curve fitting the data to the equation  $y = a + b/x$ . (data fits are good to a factor of 10%) Note that the relative levels in the concentrations of all plotted species appear to level off after 1200 seconds or about 20 minutes after the solar array was biased. It is also interesting to note that the plots in Figure 8 are equivalent to plotting partial pressure (of the indicted mass species) against the number of arcs, since the arc rate was adjusted to approximately 1 arc per second throughout the experiment.

### Summary

The presence of adsorbed water vapor at the triple junction sites on a negatively biased solar array is a necessary starting point for evidence supporting a desorbed gas molecular ionization mechanism as the main cause of arc inception. Furthermore we would expect the



**Figure 8: Quadrupole mass spectrometer results for arcs on a biased solar array.**

arcng process to drive out water vapor from the triple junction causing increasing the levels of the partial pressure of H<sub>2</sub>O over time. Ionization of the H<sub>2</sub>O by secondary electron emission and subsequent evidence of dissociative recombination processes come from observations of increased levels in the partial pressure of H<sub>2</sub>, H<sub>2</sub>O, OH and O species over time as graphically illustrated in Figure 8. Note that direct evidence of water vapor at the triple junction was confirmed from earlier spectroscopic observation of the OH molecular band in solar array arc spectra [17]. Figure 8 demonstrates that OH species ranks as the third highest concentration of partial pressure recorded in the vacuum chamber.

Dissociative recombination processes tend to occur in very short time intervals (on the order of 10 $\mu$ s). Because the acquisition time for 1 complete sweep of mass numbers 1-50 AMU (on the order of 1 minute) is much greater than the dissociative recombination time we would only expect sporadic information at best demonstrated in individual arc spectra. Even so the hydrogen peak in Figure 7a, the oxygen peak in Figure 7d and the doubly ionized oxygen O<sup>++</sup> peak in Figure 7g also begin to reveal supporting evidence of a dissociative recombination process occurring. The independent mass spectrometer

results reported tend to correlate well with earlier spectroscopic observations of the  $H_{\alpha}$  line, the  $H_{\beta}$  line and the OH molecular band lend even more weight to our hypothesis of a desorbed gas molecular ionization mechanism.

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<b>13. ABSTRACT (Maximum 200 words)</b> Previous experimental studies have hypothesized that the onset of Solar Array Arc (SAA) initiation in low-density space plasmas is caused by a desorbed gas molecular ionization mechanism. Indeed past investigations performed at the NASA Glenn Plasma Interaction Facility tend to not only support the desorbed gas molecular ionization mechanism, but have gone as far as identifying the crucial molecular species that must be present for molecular ion dominated process to occur. When electrical breakdown occurs at a triple junction site on a solar array panel, a quasi-neutral plasma cloud is ejected. Assuming the main component of the expelled plasma cloud by weight is due to water vapor, the fastest process available is due to H <sub>2</sub> O molecules and OH <sup>+</sup> ions, or more succinctly, dissociative molecular-ion dominated recombination processes: H <sub>2</sub> O <sup>+</sup> + e <sup>-</sup> → H* + OH*. Recently published spectroscopic observations of solar array arc spectra in ground tests have revealed the well-known molecular OH band (302 to 309nm), as well as the molecular SiH band (387nm peak), and the molecular CH band (432nm peak). Note that the OH band is observed in emission arcs where water vapor is present. Strong atomic lines were also observed for H <sub>β</sub> at 486nm and H <sub>α</sub> at 656.3nm in prior ground testing. Independent supporting evidence of desorbed gas molecular ionization mechanisms also come from measurements of arc current pulse widths at different capacitances. We will revisit an earlier first order approximation demonstrating the dependence of arc current pulse widths on the square root of the capacitance. The simple arc current pulse width model will be then be used to estimate the temperature of the arc plasma (currently believed to be somewhere in the range of 3 to 5 eV). The current paper then seeks to extend the outlined work by including numerous vacuum chamber measurements obtained with a quadrupole mass spectrometer. A small solar array was mounted inside the vacuum chamber. A plasma source, also mounted inside the vacuum chamber, is used to simulate a low-density plasma environment. The solar array is then biased to a high negative potential and allowed to arc while a mass spectrometer is used to record the partial pressure of H <sub>2</sub> O and to track other significant changes in mass (1 to 150) AMU.			
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